

International Journal of Heat and Mass Transfer 45 (2002) 1109-1116

International Journal of **HEAT and MASS TRANSFER** 

www.elsevier.com/locate/ijhmt

# Morphology of soot collected in microgravity droplet flames

Samuel L. Manzello a.l, Mun Y. Choi b.\*

<sup>a</sup> Department of Mechanical Engineering, University of Illinois at Chicago, Chicago, 1L 60607, USA b Department of Mechanical Engineering and Mechanics, Drexel University, Philadelphia, PA 19104, USA Received 16 March 2000: received in revised form 14 April 2001

#### Abstract

Measurements of the primary particle size. radius of gyration. fractal dimension. and the mass fractal prefactor were measured for soot sampled in n-heptane droplet flames under microgravity conditions. These represent the first suchmeasurements obtained for sphericall! symmetric droplet flames. Experiments were performed in the NASA-Glenn Research Center in Cleveland. OH. Soot was sampled at 0.2 and 0.5 s after ignition for a 2.1 mm droplet burning in atmospheric pressure air. The measured primary particle sizes were significantly larger under microgravity conditions compared to normal gravity and were found to increase with residence time. The average radius of gyrations were also larger for soot collected in microgravity conditions. Differences in the particle and agglomerate dimensions are believed to be caused by the longer residence times. The fractal dimensions were found to be nearly constant for all experiments while the prefactor term was found to increase slightly with residence time. © 2002 Elsevier Science Ltd. All rights

Keywords Soot, Fractal properties: Microgravity, Droplet combustion

## 1. Introduction

Soot is a common by-product resulting from combustion of fossil fuels. Release of soot into the atmosphere contributes to environmental and health hazards including the degradation of atmospheric visibility. ground-water contamination. and premature deaths due to ingestion. Sooting can also decrease the efficiency of energy conversion systems since carbonaceous particulates represent incomplere combustion and can cause hardware fouling. The importance of these topics has fueled numerous fundamental investigations of soot processes under normal-gravity conditions [1-5]. However. buoyancy-induced flows in normal-gravity flames produce short residence times to resolve soot formation. growth, agglomeration, and oxidation processes. Recent experiments using microgravit? environments have

'Corresponding author Tel : +1-215-895-6984; fax: +1-215-

E-mail address Choi@Drexel edu (MY Choi).

NRC-KIST Post Doctoral Fellow

produced longer residence times (by as much as three times for laminar gas-jet diffusion flames) and expanded the soot containing region to allow for accurate resolution of the soot evolution processes and for reliable soot diagnostics [6-9].

Sooting in microgravity environments present unique opportunities to investigate important mechanisms that are normally obscured under normal-gravity conditions. For example, the absence of convective flows in microgravity enables radiation as an important mode of heat transfer. Since radiation from carbonaceous particulates is much greater than from non-luminous radiation (attributed to gas species). the importance of sootin: is further amplified in reduced-gravity [I0]. Under microgravity conditions. sootin: effects that were deemed negligible, based on normal-gravity observations. were observed to be enhanced [I I]. For example. the measurements of soot concentrations for microgravity droplet flames were an order of magnitude larger than the counterpart normal-gravity flames [12] The high soot concentrations uere responsible for a greater degree of radiative heat losses and contributed to rhe reduction in the droplet burning ratr [13,14]. Investiga-

Nomenclature		K N	droplet burning rate number of primary particles
$A_a$ $A_p$	projected area of agglomerate projected area of primary particle	$rac{r_{ m c}}{r_{ m d}}$	centroid of agglomerate instantaneous droplet radius
$\frac{d_2}{D_1}$	diameter of the soot primary particle fractal dimension	$rac{r_i}{r_i}$	instantaneous flame radius centroid of pi-imary particle
do ke	initial droplet diameter mass tractal prefactor	$R_{\underline{i}}$	radius of gyration droplet burning time

tion of spherically-symmetric acetylene/air dilfusion flames supported on a porous sphere burner indicated that sooting can promote the onset of flame extinction [15]. However, despite the importance of the morphological properties of soot in the interpretation of the soot processes, there has been a lack of measurements in spherically-symmetric diffusion flames.

Characterization o i soot morphological properties including the primary particle size.  $d_{\rm p}$ , radius of gyration,  $R_{\rm g}$ , fractal dimension,  $D_{\rm f}$ , and mass fractal prefactor term.  $k_{\rm l}$  are important for the analysis of soot processes including primary particle growth, soot agglomeration and oxidation processes. Morphological measurements are *also* required for accurate formulation of optical properties (including the absorption emission and scattering constants [16]) for flame radiative emission calculations and in the interpretation of soot optical diagnostics measurements [17].

Ku et al. [7] and Konsur et al. [8] were the first to thermophoretically sample soot in microgravity laminar gas jet diffusion flames. The soot sanipled from ethylene and propane flames was analyzed using transmission electron microscopy (TEM) for primary particle size and agglomerate dimensions. The average primary particle diameter was found to be a factor of two larger under microgravity conditions compared to normal-gravity measurements. The longer residence time in microgravity enhanced the soot formation and growth processes [7]. Urban et al. [9] also recently sampled soot from microgravity laminar jet diffusion flames in experiments aboard the space shuttle. Similar trends of larger soot primary sizes for microgravity flames were observed.

Although the microgravity findings [7–9] provided important and useful information regarding the morphology of soot collected in laminar diffusion flames in microgravity, their results are not directly transferable for analysis o i soot properties of microgravity droplet combustion. There are significant differences in the residence times between the two cases. In microgravity droplet combustion, the vaporized fuel is diffused toward the flame front while undergoing pyrolysis reactions to form precursors that evolve into carbonaceous soot. Once formed near the high-temperature region, the

soot is acted upon by thermophoresis and transported towards the droplet. Opposing this mode of particle transport is the viscous drag caused by the Stefan flux (due to the vaporizing fuel droplet). The soot agglomerates ultimately reside at a radial position where these two forces are in balance, namely the sootshell location. which occurs at a non-dimensional radius  $(r_s/r_d)$  of approximately 3 Sootshells are typically observed to form within 0.1 s after ignition and linger throughout the burning lifetime. Thus, the residence time scales approximately as the square of the initial droplet diameter since the burning time  $(t_b = d_0^2/K)$  is proportional to the square of the initial droplet diameter. For a droplet of 2 mm diameter, the residence time is -6-8 sdepending on the burning rate. In contrast, the residence time for a typical microgravity laminar gas-jet diffusion flame is approximately 100 ms [9]. Spherically symmetric flames produced under microgravit) also provide a unique geometrical configuration in which to study sooting processes since the transport mechanisms of thermophoresis, diffusion, and viscous drag can be formulated using a one-dimensional model. The motivation for the present study is to sample soot within the sootshell of a n-heptane flame under microgravity conditions to investigate the influence of residence time on soot morphology. Morphological measurements including  $d_{\mathfrak{p}}$ ,  $R_{\mathfrak{p}}$ .  $D_i$ , and  $k_i$  for soot collected in microgravity droplet flames are compared to normal-gravity results.

## 2. Experimental apparatus and procedures

The microgravity experiments uere performed in the NASA-Glenn Research Center 2.2 s droptower in Cleveland, OH (see Fig. 1) Microgravity conditions are produced for 2.2 s by releasing the experimental rig into free-fall within the eight story structure. The impact of the experiment at the bottom of the tower is reduced by using a large airbag system. The experimental rig is designed to investigate the influences of sootin: on droplet combustion [14]. Fig. 2 is a schematic of the experimental rip including the 12 liter stainless steel combustion chamber. The combustion chamber contains the fuel delivery system, droplet generator, and the ignition

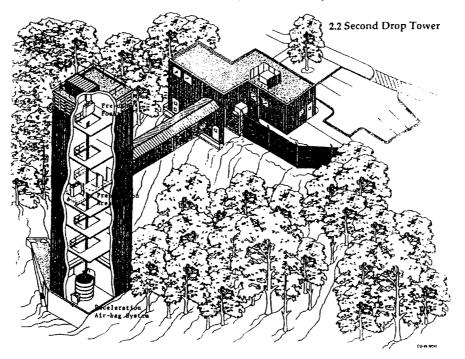


Fig. 1 Schematic of the NASA Glenn 2.2s droptower facility

assembly. The fuel droplets are generated using two opposed hypodermic needles of 0.25 mm diameter that are separated by 0.5 mm prior to the initiation of the

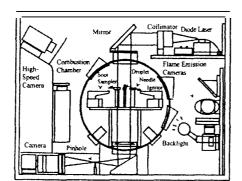


Fig. 2 Schematic of microgravity droplet combustion experimental apparatus

experiment. Fuel is pumped through the needles by solenoid-activated syringes attached to each needle. Each hypodermic needle is attached to a separate rotating galvanometric device. The exiting fuel forms a liquid bridge and the rapid rotation of the needles in opposite direction deposits the droplet onto a  $15~\mu m~SiC$  fiber. The fiber is used to fix the location of the droplet and prevent the droplet from moving out of the field of view. Approximately 0.7~s after deployment of the droplet, the fuel droplet is ignited by two hot wire igniters placed symmetrically about the droplet.

Soot agglomerates within the flame were sampled employing a thermophoretic deposition method. The sampling assembly consisted of a sinele-axis linear stepper motor (Northern Magnetics, model number 0602-2) which was controlled using an TMS motion controller (model number 1M48312). The linear stepper motor was programmed prior to each experiment for motor acceleration, velocity and sampling duration to minimize the disturbance to the droplet. A holder designed to accommodate copper TEM grids (3 mm in diameter, 300 mesh with a carbon substrate, Ted Pella, parr number 01820) was attached to the linear motor.

The entire sampling procedure was imaged using a laser bock-lit view (normally used for soot light extinction measurements) in tandem with a high speed camera operated at 200 frames per second (which was illuminated using an incandescent backlight). For the laser hack-lit view, light from 3 635 nm diode laser attached to a single tiber optic cable was expanded to a diameter of 50 mm. The beam was collimated and directed through the top optical port of the combustion chamber, focused using a 200 mm focal length plano-convex lens and reflected using a second mirror positioned at 45". The reflected beam was then imaged through a spatial filter to a high resolution CCD camera and recorded on a Betacam recorder through a fiber optic coupler.

Soot sampling was performed at 0.2 and 0.5 s after ignition for two separate microgravity experiments, to determine the influence of residence time on the soot morphology. Fig. 3 shows a schematic of the soot sampling apparatus and the microgravity droplet flame. The sampling probe insertion time through the flame was measured to be approximately 0.1 s (the rstraction time was the same). The sampling probe dwell time within the sootshell was approximately 0.2 s.

Soot agglomerates were imaged using a JEOL TEM operated at 50~K for the normal-gravity soot and at 16~ and 50~K for the microgravity soot. Fig. 1 displays images of soot agglomerates obtained from the normal-gravity flame and for the microgravity flame at 0.5~s after ignition. The micrographs obtained from the TEM analysis were placed on a light table and imaged using a high-resolution CCD camera with a 18-108~mm lens. The images were then digitized using a 480~x~512.~8~bit graylevel resolution Data Translation DT-2851~frame grabber board and a DT-2858~co-processor. The attached 18-108~com lens allowed further magnification of the image required for primary particle diameter measurements. Spatial calibration was obtained by using a separate TEM grid deposited with 100~nm polystyrene

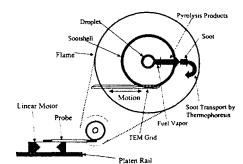


Fig. 3. Schomaiic of sampling apparatus, and of a microgravity droplet flame including the  $\,rootshell\,$ 

spheres. After the determination of a threshold (based on image enhancing technique [18]), the primary particle sizes were measured using manual edge identification [19]. The radius of gyration was measured using an automated image processing algorithm developed in the author's laboratory which is described in the next section.

#### 3. Results and discussion

The histogram distribution of the primary particle size measurements for all three experiments are displayed in Figs. 5(a)-(c). Approximately 200 panicles were analyzed for each experiment. The average primary particle size increased from 25.9 nm in normal gravity to 38.5 nm and 15.8 nm for the 0.2 and 0.5 s microgravity experiments (see Table 1). The increase in the primary particle dimensions with residence time is dramatic. For example, the volume of an average primary particle is nearly 70% larger for the  $\textbf{0.5}\,\text{s}$  case compared to the  $0.2\,\text{s}$ case. The increase in the primary particle size between normal gravity and microgravity are in qualitative agreement with soot sampled in laminar jet diffusion flames in microgravity [7.9]. At a height above burner of 10 mm. Ku et al. [7] found that the primary particle diameter of ethylene increased from LO nm in normal gravity to 38 nm in microgravity. The increase in primar), particle size 15 due to the increased residence time afforded by microgravity environments. However, there are important differences between the two geometries of the laminar gas-jet and the spherically-symmetric droplet flames. For the laminar gas-jet flames. residence time is dependent on the height above burner. Furthermore. measurements of Ku et al. [7] clearly demonstrate that there are operating conditions in microgravity which demonstrate the formation. growth and oxidation processes. These observations are not likely to occur in droplet combustion configurations since soot resides within the sootshell. far removed from the regions where oxidizers are present in appreciable concentrations [20].

Avedisian [21] performed soot primary particle measurements for toluene droplet flames in microgravity. In their experiments, soot was deposited on filter paper (placeda[ various locations on the combustion chamber) due to the impact of the experiment with the ground. Subsequent TEM analysis was performed to measure the primary particle diameters (average  $d_{\rm p}\sim 50$  nm). Since the technique neither provided temporal nor spatial information regarding the collected soot, the influence of residence times on soot morphology could not be inferred. In that study, agglomerate properties such as the radius of gyration, fractal dimension, and the mass fractal dimension were not measured.

The measurements of the radius of gyration in this study are used primarily to determine the fractal  $\ensuremath{\mathsf{I}}$ 

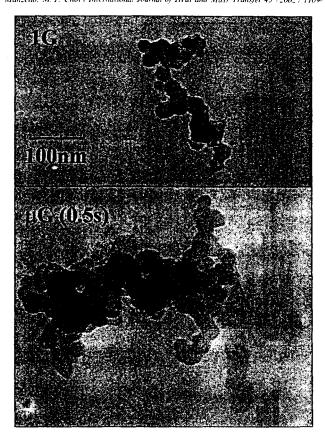


Fig. 4. Transmission electron micrograph images of soot agglomerates in normal gravity and in microgravity (0.5 s after ignition).

Table 1

Experiment	$d_{p}$ (nm) [S.D.]	$R_{y}$ (nm) [S.D.]	Range of R <sub>g</sub> (nm)	$D_{\rm f}$	k <sub>f</sub>	
Normal gravity	25.9 [4.8]	46 4 [22 7]	15-130	1.59	6.5	
ug 0.2 s After ignition	38.5[7.9]	80.0 [33.2]	28 - 204	1.56	7.0	
μg 0.5 s After ignition	45.8 [8,1]	115 6 [49 4]				

dimension and to offer insights regarding the relative differences in agglomerate sires for the three cases. In mass fractal analysis, the number of primary particles. N, comprising an agglomerate and **soot** dimensions are related by the following equation [22]:  $N = k_1 \left(\frac{R_p}{d_p}\right)^{k_1}$  (1)

$$N = k_{\rm f} \left(\frac{R_{\rm p}}{d_{\rm p}}\right)^{b_{\rm f}} \tag{1}$$

where I), is the fractal dimension and I is the prefactor term. The measurements of  $D_I$  and k; can provide important insights regarding the agglomerate growth mechanisms [4]. These parameters can be calculated based on the measurement of V,  $R_I$ , and  $d_P$ . In this study, the number of primary particles per agglomerate was calculated by measuring the projected area of the

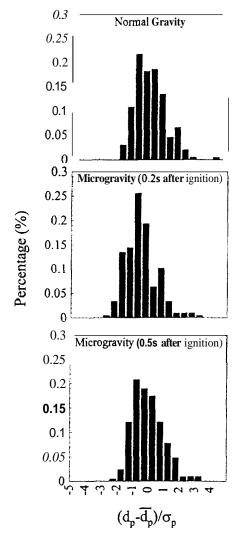


Fig. 5. Normalized histogram distributions for primary particle size measurements for normal gravity. 0.2 s microgravity, and 0.5 s microgravity experiments.

agglomerate. A., and the projected area of the pi-imary panicle, 4, (based on mean  $d_{\rm p}$  measurements)

$$N = 1.0 \left(\frac{A_a}{A_p}\right)^{1.08}$$

The radius of gyration based on the projected image of the agglomerate was measured using a custom-image processing algorithm. The radius of gyration of an agglomerate is defined by the following relationship:

$$R_{\rm g} = \sqrt{\sum_{\rm rel}^{N} \frac{(r_{\rm t} - r_{\rm e})^2}{N}}.$$
 (3)

where  $r_i$  is the centroid of the primary particle and  $r_c$  is the centroid of the agglomerate. Since it is difficult to identify each individual particle that constitute an agglomerate (due to the sheer number of particles concerned which can range from a few to a few thousands), the  $R_x$  calculation formulation was revised by considering each pixel to be the individual element of interest [22]. The equation for the radius of gyration remains the same but N now corresponds to the total number of pixel elements comprising the agglomerate (typically on the order of 20-30 thousand pixels in our experiments).

The average  $R_s$  for the normal-gravity experiment was 46.4nm (based on 92 agglomerates) whereas to: the 0.2 and 0.5 s microgravity experiments, values of 80.0 nm (146 agglomerates) and 115.6 nm (119 agglomerates) were measured. respectively. This increase is a manifestation of the exrended residence time (due to influence of thermophoresis in 'trapping soot' within the sootshell) during which additional growth and agglomeration can occur. The differences among the experiments are not only apparent uhen comparing the mean values but also in the ranges of agglomerate sizes measured The ranges in R<sub>s</sub> for the normal gravity. **0.2** and 0.5 s cases are 15-130, 28-204, 37-270 nm. respectively. Since the transport of soot agglomerates formed near the high-temperature region towards the sootshell occurs throughout the burning lifetime. the agglomerates within the sootshell should display wider distribution of newly formed and mature soot of varying age as residence time increases.

The measurements of N,  $R_8$  for each agglomerate and the average value of  $d_p$  for each condition were used to calculate  $D_f$  and  $k_f$ . As shown in Figs. 6(a)-(c) (for measurements obtained in normal gravity. 0.2 s microgravity, and 0.5 s microgravity), the fractal dimension represents the slope of the ln[N] versus  $ln[R_g/d_p]$ and the prefactor term is the intercept at  $R_s = d_p$ . These results are also summarized in Table I. Unlike the physical dimension measurements, D<sub>i</sub> was nearly constant for all three cases.  $D_{\rm f}$  for normal-gravity soot was 1.59, whereas for the 0.2 and 0.5 s, cases they were 1.56 and 161, respectively. The fractal dimensions lor the three cases are consistent with previous normal-gravity results obtained tor gaseous (acetylene, propylene, ethylene, propane) turbulent diffusion flames [23]. In that stud:. it was found that the average  $D_f$  (measured using similar approach as described in the present study) was approximately 1.65 [23]. The investigation of Ku et al

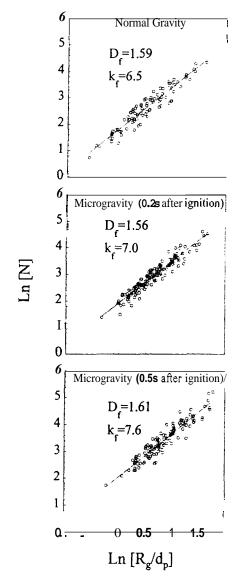


Fig. 6 Plot of  $\ln(N)$  versus  $\ln(R_{\rm g}/d_{\rm p})$  for soot measurements in normal gravity, 0.2 s microgravity, and 0.5 s microgravity experiments

[7] is the only other study for which fractal dimensions were measured for soot produced in microgravity diffusion flames. Ku et al. [7] measured the fractal dimen-

sion for ethylene and propane that ranged from 1.41 to 1.93.

The prefactor term. in addition to the fractal dimension. is essential for the accurate characterization of the soot morphology and the agglomeration process [24] The magnitude of the prefactor tenii represents the compactness, i.e., the level of packing of the soot priman particles that comprise the agglomerate [25]. In our esperinients, the average  $k_{\rm f}$  increased slightly from 6.5 in normal gravity to 7.0 and 7.6 for the 0.2 and  $0.5\,s$ microgravity experiments, respectively. The increase in the number of primary particles comprising an agglomerate scales linearly with  $k_f$ . There are no other measurements of  $k_i$  (based on  $R_r$ ) for soot collected in microgravity flames. The present values are. however. within the reported measurements of  $k_1$  for normal- ${\it grav}$  ity experiments which range from approximately 4.0 for soot collected in a laminar premixed flame [24] to 8.5  $\ensuremath{\text{for}}$  soot collected in the overfire region of tui-bulent diffusion flames [23]. Additional esperiments will be required to fully investigate this behavior.

These experiments have demonstrated that the morphological properties of soot produced in microgravity droplet flames are significantly different than those measured in normal-gravity experiments. Furthermore, the spherically-symmetric droplet configuration is unique in that the residence time can be controlled over a range beyond what can he obtained for laminar gas-jet diffusion flames. Additional experiments using facilities producing longer microgravity observation times are required to investigate whether the observed trendi of the soot physical and fractal dimensions continue with increasing residence times.

# Acknowledgements

The authors gratefully acknowledge support through NASA (NCC3-655 and NCC3-822. Dr. Paul Ferkul serving as project scientist). S.L.M. also acknowledges support from a NASA GSRP fellowship. The assistance of the NASA-GRC staff (Messrs. A Birchenough. M. Johnston. R. Mileto. C. Hampton. J Carrion and J. Owens) is sincerely appreciated.

# References

- R.J. Santoro, T. Yeh, T.T. Horvath, The transport and growth of soot particles in laminar diffusion flames, Combust. Sci. Technol, 53 (1987) 89–115
- [2] S.J. Harris, A.M. Weiner, Surface growth of soot particles in premixed ethylene/air flames, Combust Sci Technol 31 (1983) 155–167
- [3] A. D'Alessio. Laser light scattering and fluorescence diagnostics of rich flames produced by gaseous and liquid

- fuels, in. D. Siegla, G. Smith (Eds.), Particulate Carbon. Formation during Combustion. Plenum Press, New York, 1981.
- [4] C.M. Megaridis, R.A. Dobbins, Morphological description of flame-generated materials. Combust. Sci. Technol. 71 (1990) 95–109.
- [5] J. Zhang, C.M. Megaridis, Soot microstructure in steady and flickering laminar methane/air diffusion flames. Combust. Flame. 112 (1998) 373–484.
- [6] C.K. Law, G.M. Faeth, Opportunities and challenges of combustion in microgravity, Prog. Energy Combust. Sci. 20 (1994) 65-113.
- [7] J C Ku, DW Griffin, PS Greenberg, J Roma, Buoyancy-induced differences in soot morphology. Combust Flame 102 (1995) 216–218.
- [8] B Konsur, C.M Megaridis, D Griffin, Fuelpreheateffects on soot-field structure in laminar gas jet diifusion flames burning in 0-G and I-G. Combust. Flame 116 (1999) 334-337.
- [9] D.L. Urban, Z.G. Yuan. P.B. Sunderland. G.T. Linteris. J E Voss, K.C. Lin, Z. Dai. K. Sun. G.M. Faeth. Structure and soot properties of non-buoyant ethylene/air laminar jet diffusion flames. AIAA J (1998) 1346-1360.
- [10] R Siegel, J. Howell. Thermal Radiation Heat Transfer. third rd., Hemisphere, New York, 1981
- [11] B.D. Shaw, F.L. Dryer, F.L. Williams, J.B. Haggard Jr., Sooting and disruption in spherically-symmetrical combustion of decane in air. Acta Astron. 17 (1988) 1195–1202.
- [12] M Y. Choi, K.O. Lee Investigation of sooting in microgravity droplet combustion. Proc Combust. Inst. 26 (1996) 1149–1153.
- [13] G.S. Jackson, C.T. Avedisian, The effect of initial droplet diameter in sphericall) symmetric droplet combustion of sooting fuels. Proc. R. Soc. London A 446 (1994) 255–276.
- [14] K.O. Lee, S.L. Manzello, M.Y. Choi. The effects of initial diameter on sooting and burning behavior of isolated droplets under microgravity conditions, Combust Sci Technol 132 (1998) 139–156

- [15] S. Berhan, A. Atreya, D. Everest, Radiant extinction of gaseous diffusion flames, in: Proceedings of the Fifth International Microgravity Combustion Workshop, NASA/CP 20017, Charlend OH 1999, pp. 43-46.
- NASA/CP 208917, Cleveland, OH. 1999, pp. 43-46.
  [16] R.A. Dobbins, G.W. Mulholland, N.P. Bryner, Comparison of a fractal smoke optics model with light extinction measurements. Atmos. Environ. 26 (1994) 873-887.
- [17] U.O. Köylü, G., M.-Faeth, Structure of overfire soot in buoyant turbulent diffusion flames at long residence times. Combust. Flame 89 (1992) 140-156
- [18] M.Y. Choi. Droplet combustion characteristics under microgravity and normal gravity conditions. Ph.D. Thesis. Princeton University. 1992
- [19] J. Zhang. The influence of ferrocene additive on soot formation in laminar ethylene diffusion names. Ph.D. Thesis. University of Illinois. Chicago. 1995.
- [20] A.J. Marchese, F.L. Dryer, V. Nayagam, Numerical modeling of isolated n-alkane droplet flames: initial comparisons with ground and space-based microgravity experiments. Combust. Flame 116 (1999) 432-459.
- [21] C.T. Avedisian. Soot formation in spherically symmetric droplet combustion. in F.L. Dryer. R.F. Sawyer (Eds.). Physical and Chemical Aspects of Combustion. Gordon and Breach. New York. 1996. pp 135–160.
   [22] R. Jullien, R. Botet. Aggregation and Fractal Aggregates.
- [22] R Jullien, R. Botet. Aggregation and Fractal Aggregates World Scientific, New York. 1987.
- [23] Ü.Ö. Köylü, G.M. Faeth. T.L. Farias, M.G. Carvalho, Fractal and projected structure properties of soot aggregates, Combust. Flame 100 (1995) 621-633.
- [24] I Cai. N. Lu. C. Sorenssn. Analysis of fractal cluster morphology parameters: structural coefficient and density autocorrelation function cutoff, J Colloid Interface Sci. 171 (1995) 470–473.
- [25] S Wu, S.K. Fncdlander, Note on the power law equation for fractal-like aerosol agglomerates. J. Colloid Interface Sci. 159 (1993) 246–248